

# Definition of longitudinal emittance

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$$\varepsilon_{\Delta\phi\Delta w} = 2\pi W_p / \lambda \varepsilon_{nz}$$

$$\varepsilon_{\Delta\phi\Delta w} (\text{MeV.deg}) = 3.652\text{E}5 \varepsilon_{nz} (\text{m})$$

0.13 mm.mrad

0.364E-6

# Required DTL focusing gradient

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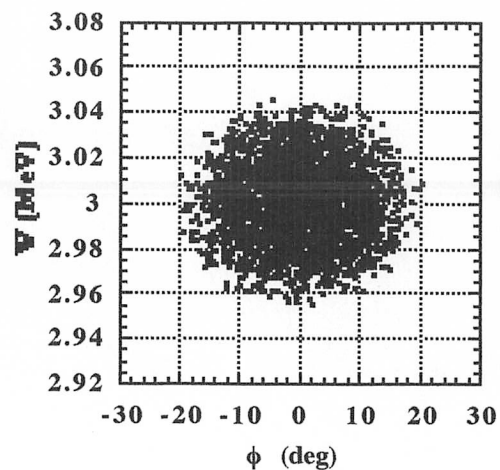
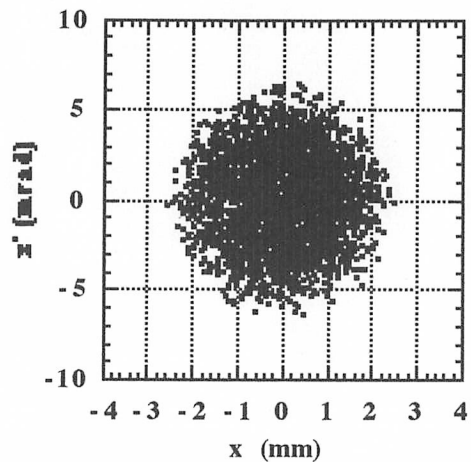
$$\mathbf{B}' \propto \frac{\mathbf{Freq}^2}{\beta} \quad (\beta = v/c)$$

- **make an RFQ of higher output energy**

**In JHF design, it was solved by  
324 MHz, 3-MeV RFQ and FD focusing  
scheme in DTL, resulting in some decrease in  
ZTT in DTL due to rather large DT diameter**

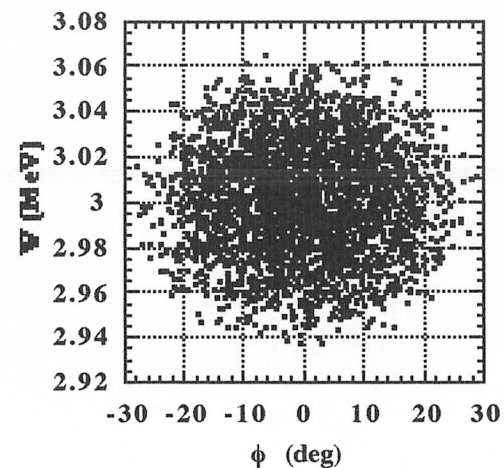
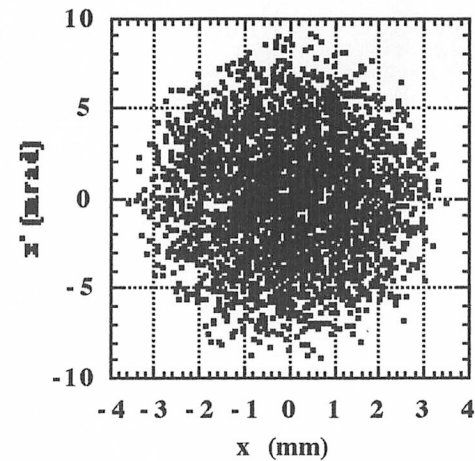
# Type A and Type B input beams

## Type A



## Type B

B



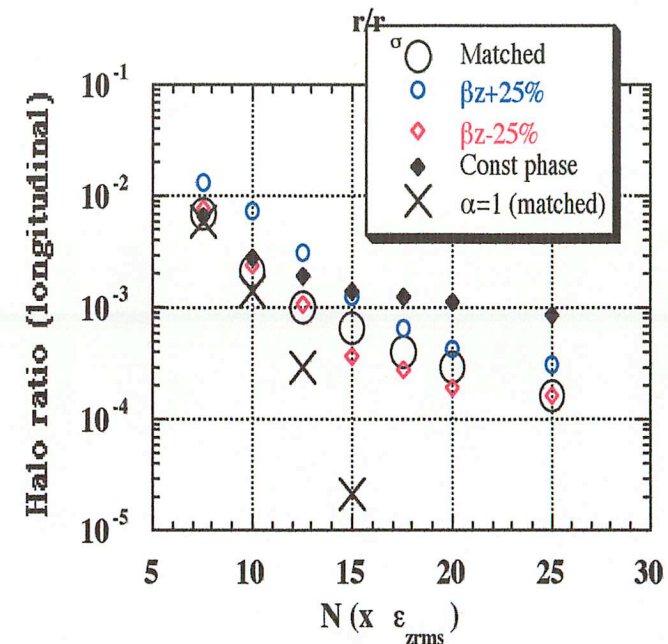
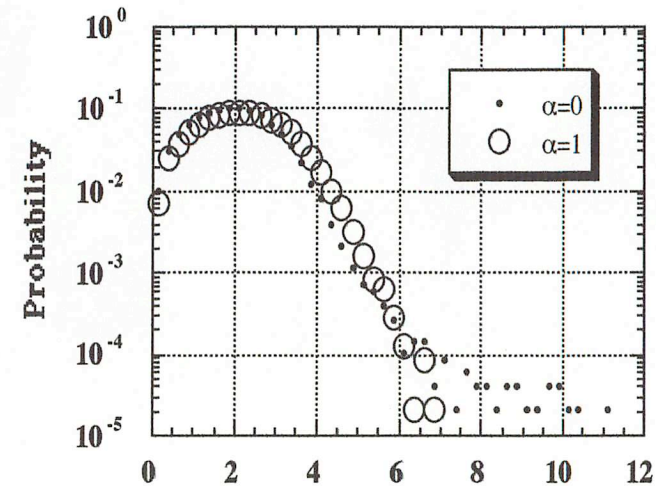
# Definition of halos

## Transverse

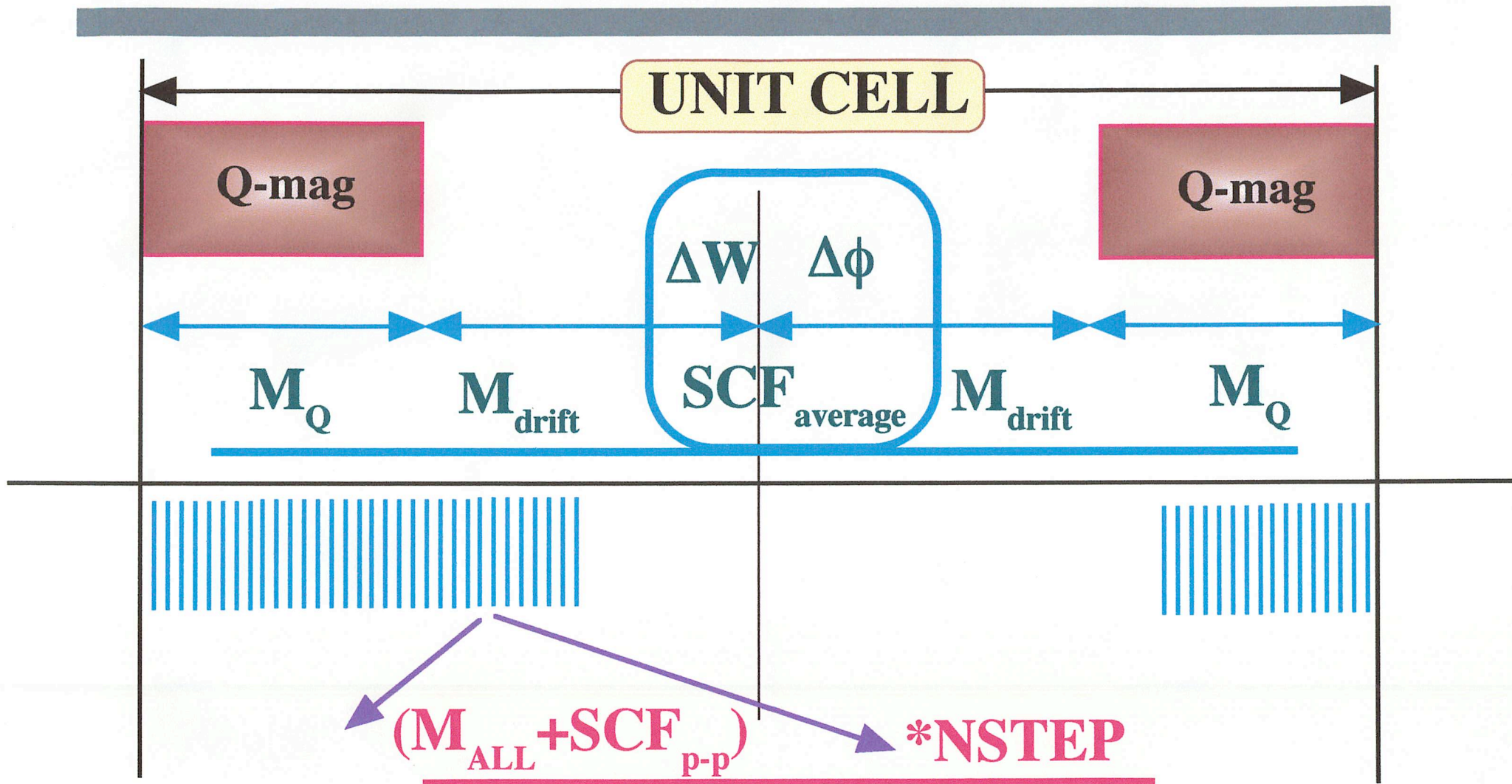
Halo-like particles are defined by those in the outside of 6.5 times the standard deviation of the radial distribution of the output beam.

## Longitudinal

Halo-like particles are defined by those in the outside of 12.5 times the longitudinal output rms emittance.



# Details in calculation (PARMILA vs. LINSAC)



# Is emittance growth in LINSAC large?

Constant phase advance of 60 deg, 30 mA, 3 - 148 MeV

Type A input beam  $\varepsilon_x=0.187$ ,  $\varepsilon_z=0.364$

transverse emittance growth with  $\alpha=1$  ---->  $\sim 10\%$

transverse emittance growth with  $\alpha=0$  ---->  $\sim 20\%$

- field energy
- nonlinear space-charge effects
- injection mismatch
- something arising from the code
  - true
  - insufficient approximation
  - error

# Emittance ratio and emittance growth

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<b>Beam</b>	$\epsilon_{nx}$	$\epsilon_{nz}$	$\epsilon_{nz}/\epsilon_{nx}$	<b>EGF-x</b>	<b>EGF-z</b>
<b>Type A</b>	<b>0.187</b>	<b>0.364</b>	<b>2</b>	<b>1.4</b>	<b>1.3</b>
<b>Type B</b>	<b>0.372</b>	<b>0.728</b>	<b>2</b>	<b>1.3</b>	<b>1.1</b>
<b>4L-beam</b>	<b>0.744</b>	<b>0.364</b>	<b>0.5</b>	<b>1.2</b>	<b>1.4</b>

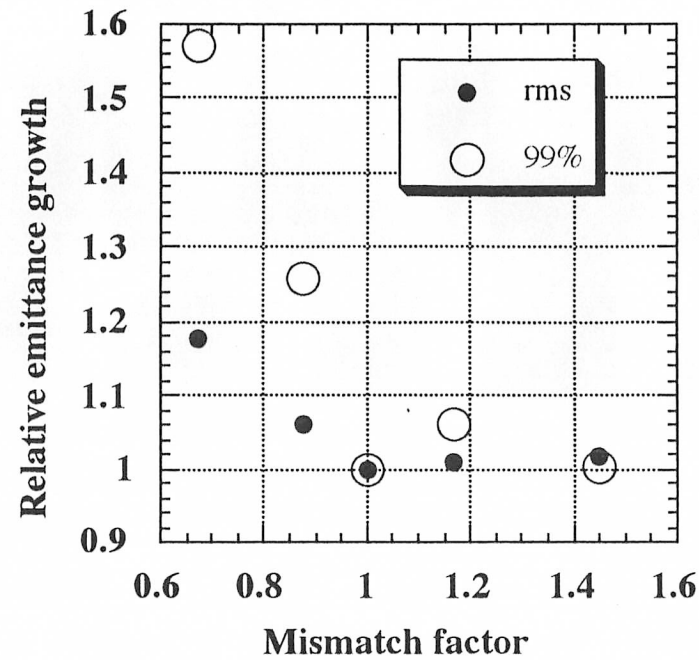
# $\alpha$ -dependence of emittance growth

<b>Macro simulation</b>		
Transverse	1.53	( $\alpha=0$ )
	1.43	( $\alpha=1$ )
<b>Real simulation</b>		
Transverse	1.46	( $\alpha=0$ )
	1.40	( $\alpha=1$ )

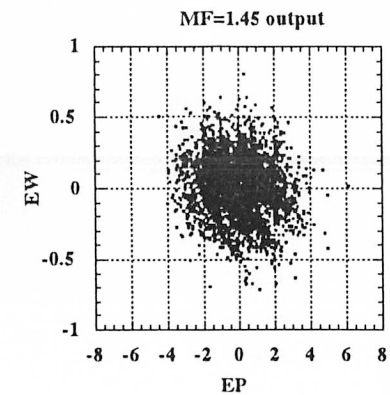
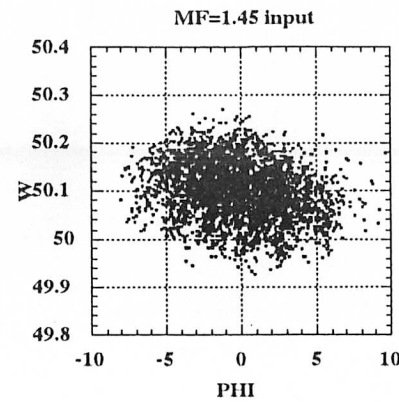
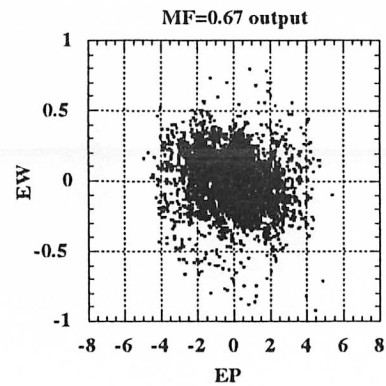
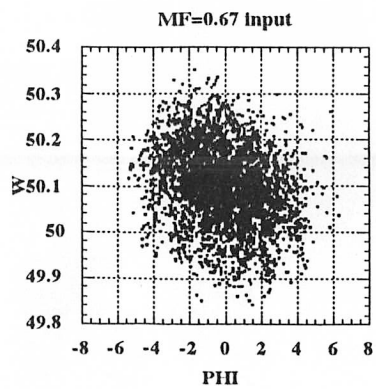
Although there are large differences in the Q-value for the macro and real simulations ( $Q \sim 12000$  and  $Q=1$ ), the  $\alpha$ -dependence of the change in the transverse emittance is nearly the same for the macro and real simulations. Therefore, it is concluded that the effects of large Q-value in the macro simulation upon the growth of rms transverse emittance is very small. A similar situation can be expected in the growth of the longitudinal emittance, suggesting that the large longitudinal emittance growth obtained in the constant phase advance simulations is not due to the effects of large Q-value, although the effects of  $\alpha$  on the longitudinal emittance growth is very large.



# Longitudinal matching in HEBT (2)

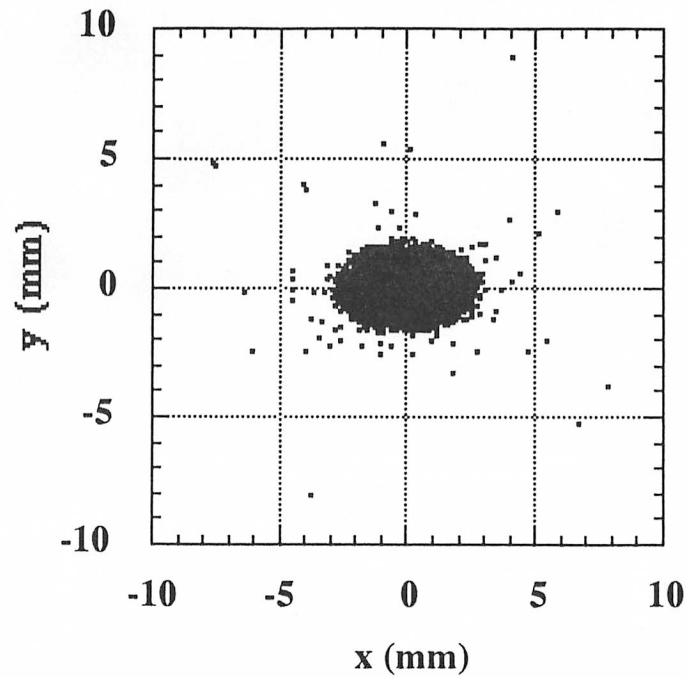


Type A beam is used.  
MF=0.86 is DTL output.

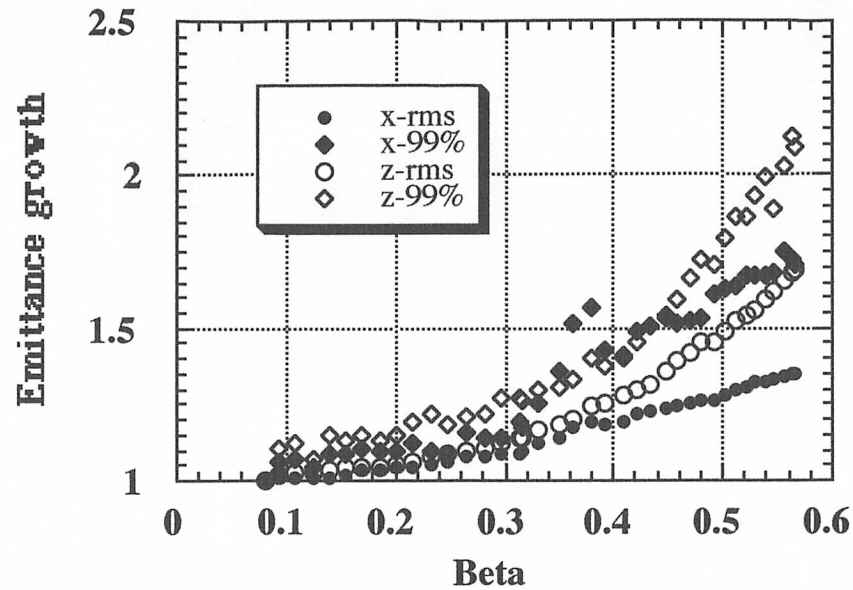


# Transverse halos in constant phase advance

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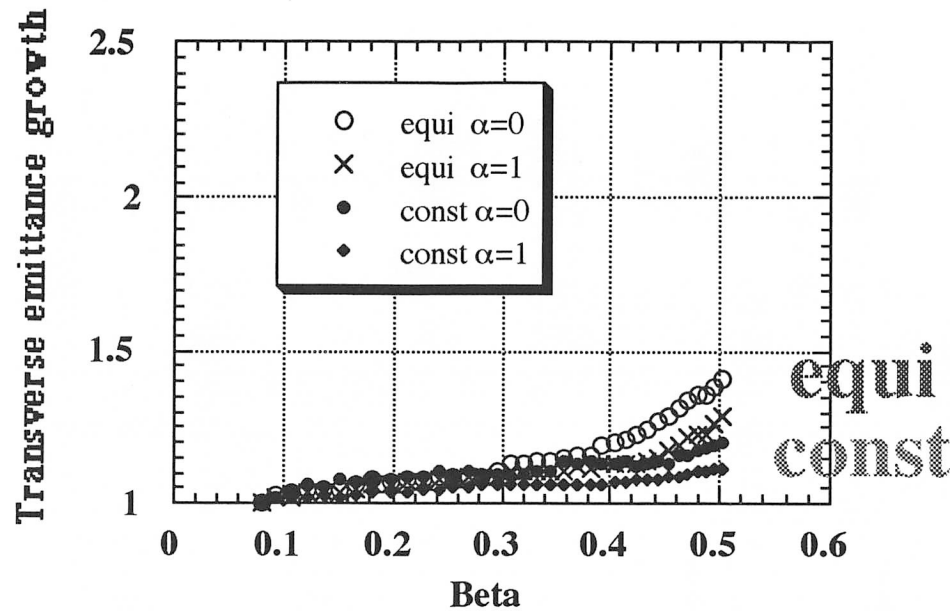


# DTL + SDTL (30-mA simulation)

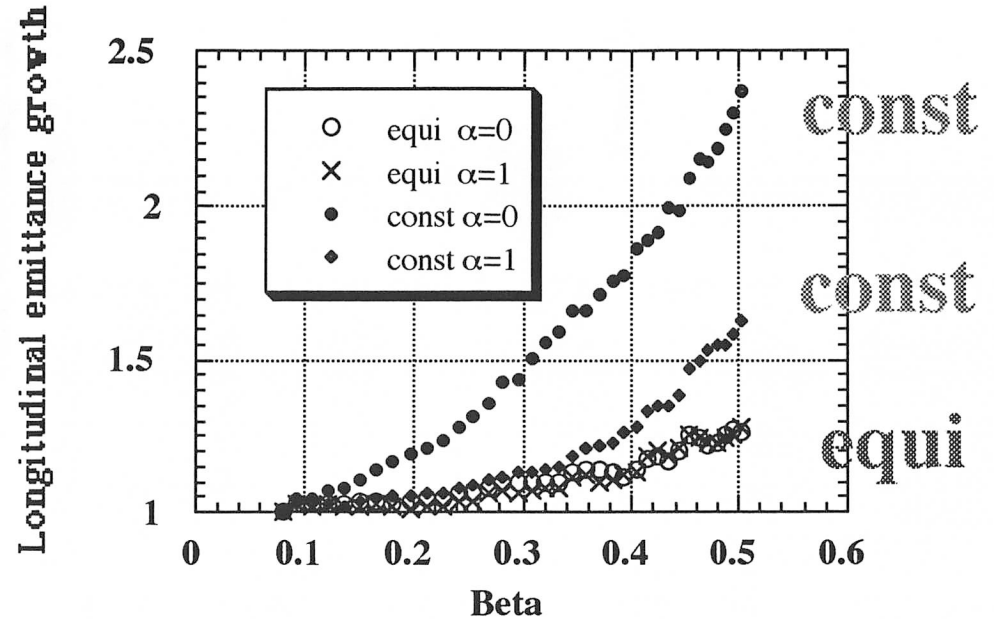


**Fig. 4.25** Emittance growth (rms and 99%) along the linac. A 30-mA type-B injection beam is used.

# Emittance growth for both focusing methods



Transverse



Longitudinal

## LINSAC results

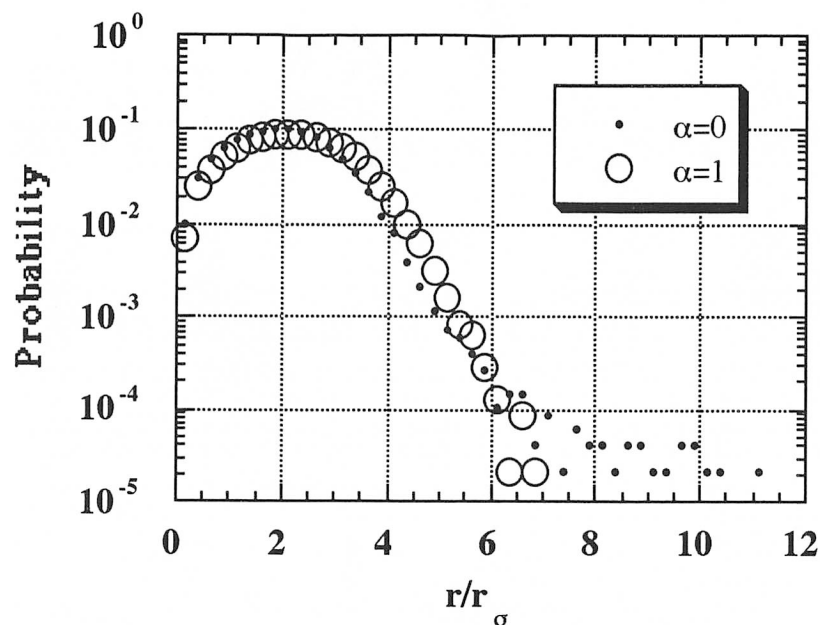
DTL acceleration from 3 to 148 MeV, 30-mA beam

$\alpha=0$  means with collision effects

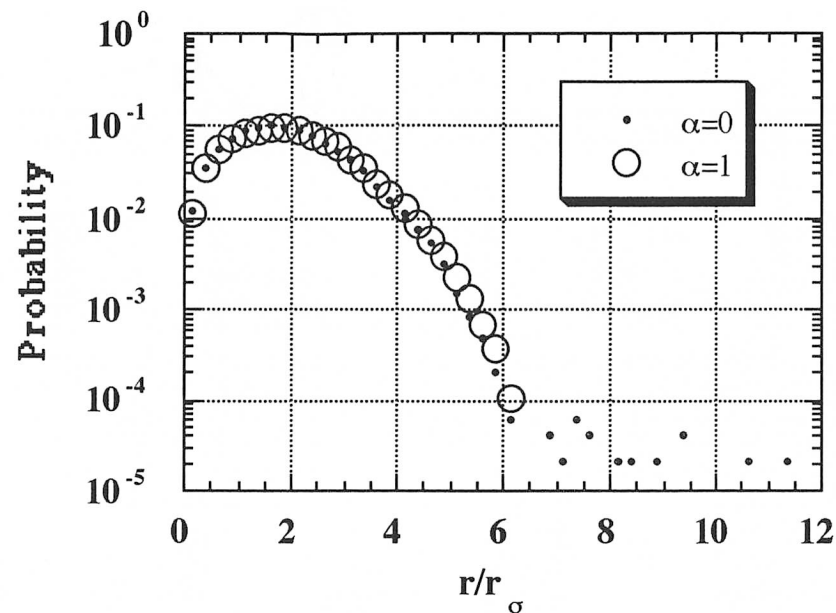
$\alpha=1$  means without collision effects

# Radial distribution (macro and real sim.)

MACRO



REAL



Comparison of the radial distribution for the matched macro simulations of  $\alpha=0$  and  $\alpha=1$ . The radii were normalized by the standard deviation ( $r_\sigma$ ) for each profile. The number of particles is 48000.

Comparison of the radial distribution for the matched real simulations of  $\alpha=0$  and  $\alpha=1$ . The number of particles is 48000.  $Q=1$ .

# Acceptance

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**Acceptance of the DTL  
(normalized, 90%)**

**x: 43  $\pi$ mm-mrad**

**z: 9.3  $\pi$ MeV-deg**

**Acceptance of the SDTL  
(normalized, 90%)**

**x: 21  $\pi$ mm-mrad**

**z: 40  $\pi$ MeV-deg**

# Equi. vs Const phase

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	<b>rms</b>	<b>90%</b>	<b>99%</b>	<b>99.9%</b>
<b>EQUI</b>	<b>0.25263E-01</b>	<b>0.10969E+00</b>	<b>0.20726E+00</b>	<b>0.41796E+00</b>
<b>CONST</b>	<b>0.23442E-01</b>	<b>0.10052E+00</b>	<b>0.21223E+00</b>	<b>0.43362E+00</b>

# Reducing beam losses in the linac

- less increase in emittance growth with higher frequency
- equal increase in both emittance growth
- less halo formation by equipartitioning focusing
- keep ratio between transverse and longitudinal focusing strength through the linac
- avoid longitudinal transition up to 200 MeV
  - In upgrade, two or three times the frequency for CCL instead of more than four
- sufficiently large bore radii and acceptances
- Stabilized rf structures and short SDTL unit tank
- Simple DTL injection scheme, avoiding tuning errors
- small repetition frequency, reducing transient loss



# Intrabeam scattering

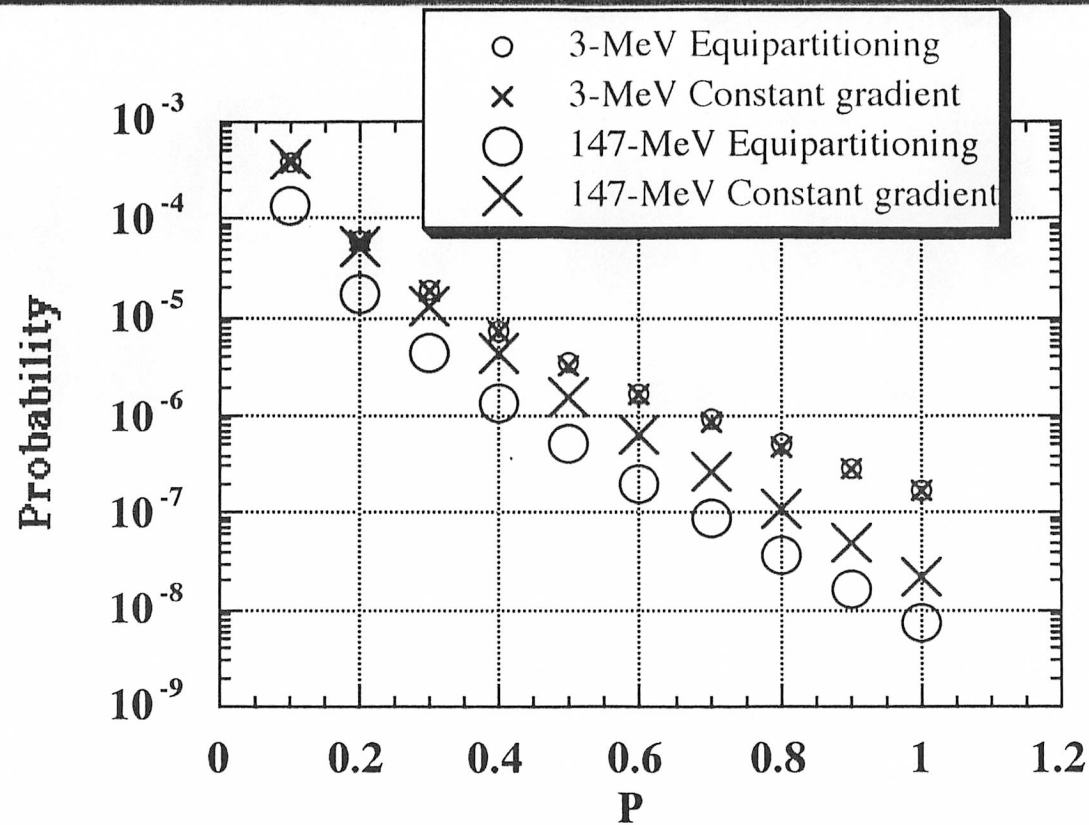
The rate of collisions due to particles having a relative longitudinal velocity of more than  $v_{th}$  is given by

$$\left( \frac{dN}{dt} \right)_{collision} = \frac{N^2 r_c^2 c \Lambda^3}{16\pi \sigma_x \sigma_y \sigma_z} D(\xi)$$

$$D(\xi) = \sqrt{\xi} \left\{ -\frac{3}{2} e^{-\xi} + \frac{\xi}{2} \int_{\xi}^{\infty} \frac{\ln(u) e^{-u}}{u} du + \frac{1}{2} \left( \xi - \xi \ln(\xi) + 2 \right) \int_{\xi}^{\infty} \frac{e^{-u}}{u} du \right\}$$

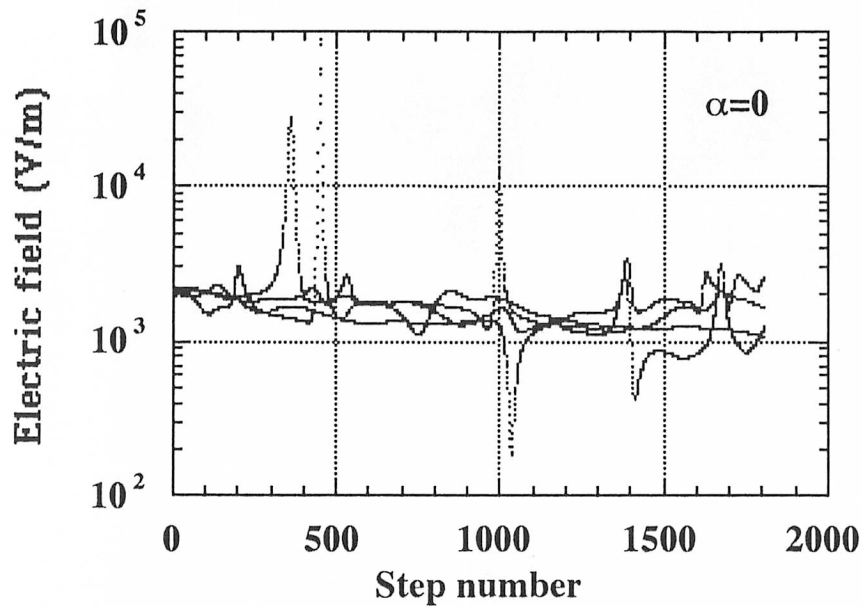
where  $N$  is the number of particles in a bunch;  $r_c = 1.53 \times 10^{-18}$  m the classical proton radius;  $\Lambda = c/v_{th}$ ;  $x = v_{th}/v_{rms}$ ;  $v_{rms}$  the rms longitudinal velocity in the moving frame; and  $\sigma_x$ ,  $\sigma_y$  and  $\sigma_z$  the rms radii of the bunch for  $x$ ,  $y$  and  $z$  directions, respectively.

# Collision probability

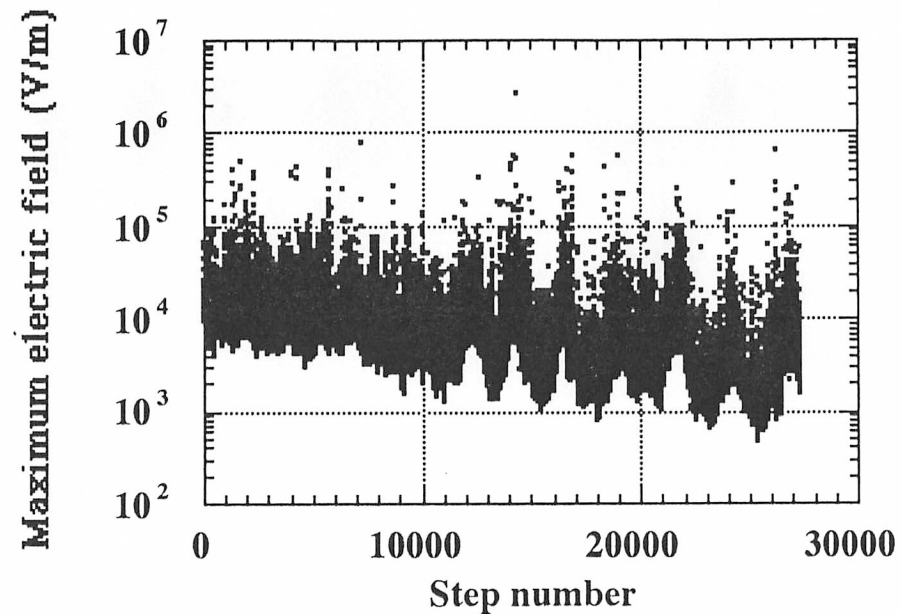


Collision probabilities at two kinds of energies ( 3 and 147 MeV) for two kinds of focusing methods (the equipartitioning method and the constant gradient method). P expresses the ratio of the relative longitudinal velocity to the threshold longitudinal velocity, which is set to the average of the rms transverse velocities. The probabilities stand for the period for the transit time through the linac.

# Collision in the real simulation - Coulmb forces

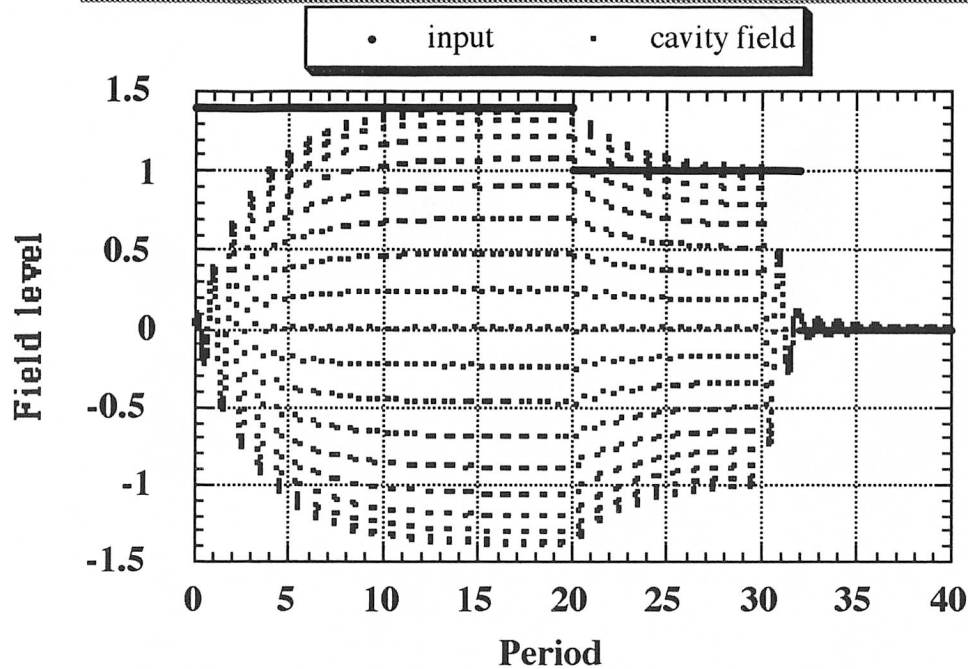


Electric field on four particles among 1600 particles.

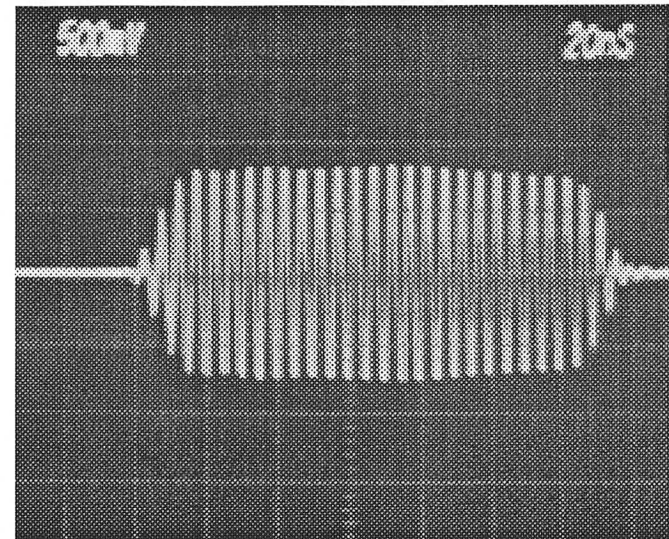


Maximum electric fields due to the space-charge effects among all particles in every calculation step up to 27324 for the real simulations.

# RF chopper - waveform



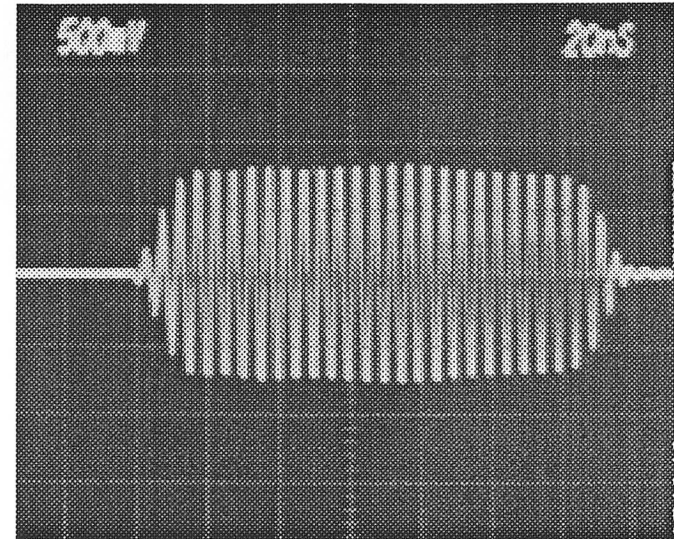
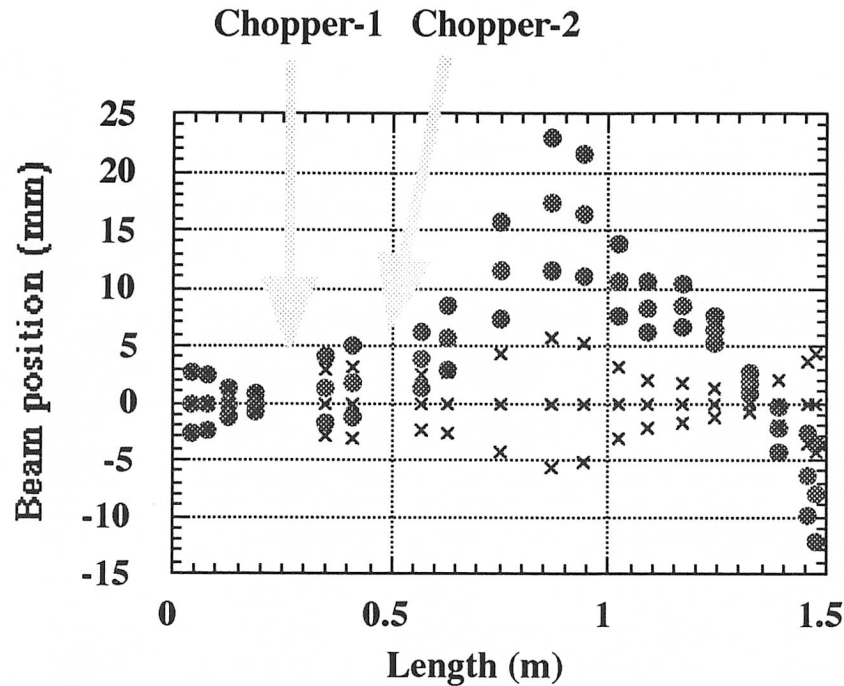
**8.5 kW peak power**



**Results of rf-field  
simulation in a low-Q  
cavity**

In order to reduce losses during rise and falling time,  
modulation in the ion source is desirable.

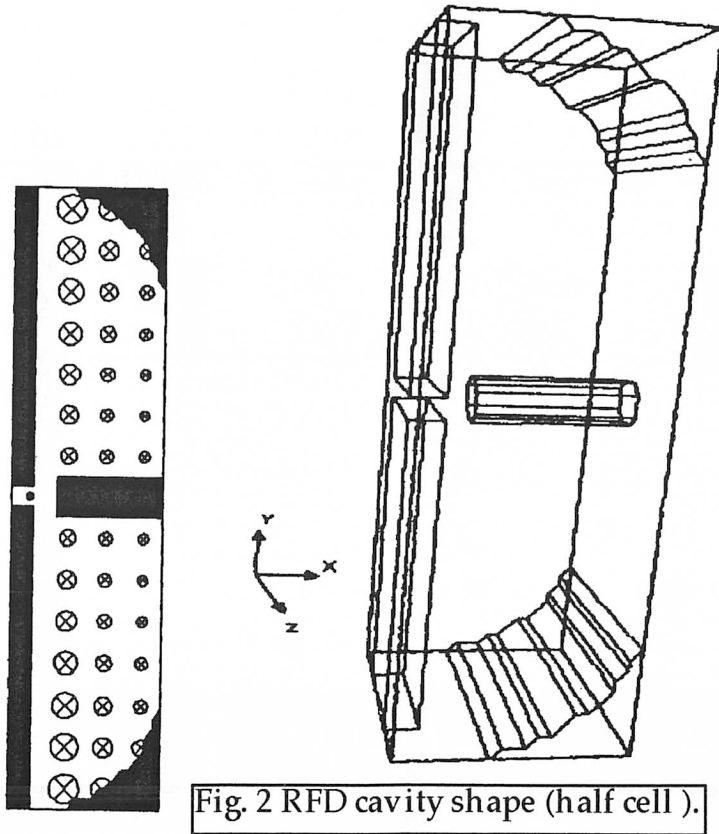
# RF chopper



Traces of the chopped (circles) and normal (cross) beams along the MEBT. The symbols indicate the outermost points and the center of the beam. Each chopper gives 15mrad defelection.

Shape of an rf pulse for the chopper cavity, measured in the preliminary experiment at a frequency of 201 MHz. A peak power is about 8.5 kW. 20nsec/div.

# Chopper cavity



Required rf parameters for the chopper system (432 MHz).

Number of cavities	Deflection mrad	Pc/cavity kW	Total power kW
1	14	58.2	58.2
2	6	11.0	22.1
4	4	7.1	28.5